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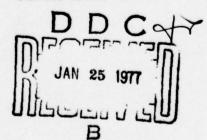
CALCULATIONS OF PRESSURE COEFFICIENTS AND FORCES FOR THE USNS WYMAN WITH AN ADDED STRUT AND FOIL CONFIGURATION

by

N.M. White and H.T. Wang

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SHIP PERFORMANCE DEPARTMENT
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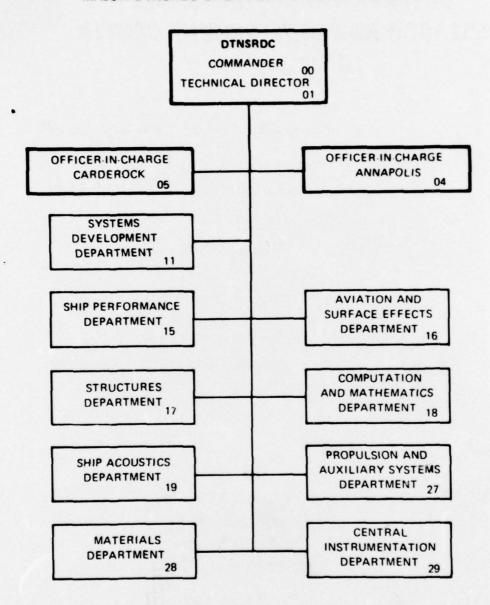


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ABSTRACT

A selected foil and strut configuration to be added to the hull of the USNS WYMAN was subjected to computer analysis to obtain the forces and pressure distributions over the foil. Both two-dimensional and three-dimensional models were developed and compared to obtain both a range of predicted lift coefficients for the given configuration, and the effect of the presence of a strut on the foil. The presence of a strut and non-zero foil pitch angle were both found to significantly depress the pressure coefficients over the foil surface as well as increasing the total applied force on the foil.

ADMINISTRATIVE INFORMATION

This work was sponsored by the Naval Ship Engineering Center under Work Procurement Request #WR-N65-197-76-WR of 30 June 1976 with internal job order number 1-1552-150-01.

INTRODUCTION

NAVSEC Code 6136 requested DTNSRDC to provide calculations of pressure distributions and lift coefficients on a selected configuration of foil beneath the hull of the USNS WYMAN. The results of calculations are given here in several forms in answer to that request. In order to provide as thorough a study as possible, calculations were made for both three-dimensional and two-dimensional models of the ship-strut-foil configuration illustrated in Figure 1A.

Two foil orientations were considered. In the first instance (illustrated in Figure 1B, CASE 1), the foil is parallel to the ship bottom, and the onset flow is at an angle, β , relative to both the ship bottom and the foil chord. As seen in Figure 1B, CASE 1, a positive angle of onset flow, β , corresponds to the case where the vertical component of the flow is directed upward toward the ship bottom. In this case, the ship plus foil may be viewed as undergoing pitch and/or heave motions in parallel flow. In the second instance (shown in Figure 1B, CASE 2), the foil has a pitch angle, α , relative to the ship bottom, while the onset flow remains parallel. Figure 1B, CASE 2, defines a positive value of α as a nose-up pitch of the foil relative to the ship's bottom. The dimensions shown in Figure 1A, as well as the illustrated coordinate system and conventions for angle definition, are those used throughout this study.

PRESENTATION OF RESULTS

THREE-DIMENSIONAL FOIL

Pressure coefficients for the three-dimensional configurations were computed by using the Douglas lifting body program, described in Reference (1). The following three cases were calculated: (1) the finite aspect-ratio foil having NACA 0020 sections deeply submerged ($h\rightarrow\infty$ in Figure 1B, CASE 1) in the absence of the ship or strut at onset flow angles, β , of 0° and $+2^{\circ}$; (2) the same foil with pitch angles, α , of 0° and $+2^{\circ}$ placed 2 feet (0.61 m) below a horizontal flat plate in parallel flow (Figure 1B, CASE 2), and (3) the ship-strut-foil configuration in its entirety at angles of onset flow, β , equal to 0° , $\pm 2^{\circ}$, and $\pm 5^{\circ}$ simulating pitch or heave of the entire configuration (Figure 1B, CASE 1).

The computed pressure coefficient distributions for various portions of the foil are shown in Figures 2 through 9. These calculations do not include the effect of the ship boundary-layer displacement thickness, δ^* , on the pressure distributions. The value of δ^* was estimated to be 0.1 feet (0.03 m) which would reduce the clearance between the ship and foil. However, since the actual gap between the ship bottom and foil is everywhere at least two feet (0.61 m), it was assumed that the displacement thickness could be neglected.

The Appendix contains illustrations showing the extent of geometric detail employed in the numerical modeling of the various components of the ship-strut-foil configuration. The modeling was done using digitized data taken directly from the ship's lines. A fine grid was used on the foil and strut sections, with data lines taken approximately 1.3-1.5 feet apart (0.40-0.46 m). A coarser grid was used on the ship since its shape is simple, with an average quadrilateral being about 1.3 feet wide and 26.5 feet long (0.40 x 8.08 m). At the intersection of the strut with the hull and the strut with the foil, quadrilaterals were spaced about every 0.5 ft (0.15 m) in each direction, allowing for more accurate calculations at the areas of maximum interest.

Information on the pressure coefficients, C_p , on the foil is supplied here in five forms. Figures 2-6 show the chordwise distributions of C_p for the three-dimensional equivalent of Figure 1B, CASE 1, at three spanwise foil locations relative to the strut for five onset flow angles β . Secondly, Figure 7 shows the spanwise distribution of the minimum values of C_p on the upper and lower surfaces. In addition, the changing chordwise location of the line of minimum C_p is also indicated. Thirdly, Figures 8 and 9 respectively provide a comparison of the chordwise distributions of C_p at angles α of 0^0 and $+2^0$ for the deeply submerged foil and the foil placed two feet below a flat plate. In addition, Figure 10 illustrates changes in the local lift coefficient, C_{LL} , across the foil span. Lastly, Table 1 gives the lift coefficient, C_L , which is the integrated effect of the pressure coefficients, for the three cases at the various angles considered. The lift coefficient is given by

$$C_{L} = \frac{L}{(1/2)\rho U^{2}C(2S)}$$
 [1]

where

L is the lift on the entire foil

p is the fluid density

U is the advance speed of the ship

C is the chord = 10 feet (3.05 m)

S is the semispan = 15 feet (4.57 m)

TWO-DIMENSIONAL FOIL

Since the three-dimensional lifting body program has seldom been used at DTNSRDC, it was decided to make some calculations on two-dimensional configurations to check its accuracy. These calculations were made by using the potential flow computer program XYPF, which is based on the theory given in Reference (2) with corrections made to provide for lift.

Modifications were also made to allow for a horizontal plane of symmetry. The following two cases were calculated: (1) the NACA 0020 section deeply-submerged at foil pitch angles, α , of 0° , $\pm 2^{\circ}$, and $\pm 5^{\circ}$, and (2) the same section placed two feet below a flat plate with foil pitch angles, α , of 0° , $\pm 2^{\circ}$, and $\pm 5^{\circ}$.

Figure 11 shows the pressure coefficients for the case of the section submerged below a flat plate at pitch angles, α , of 0° , $\pm 2^{\circ}$, and $\pm 5^{\circ}$. The pressure coefficients for the deeply submerged foil at pitch angles, α , of 0° and 2° may be found in Figures 12 and 13 respectively. For comparison, also included in Figures 12 and 13 are the equivalent cases of the two- and three-dimensional foils submerged below a flat plate, the deeply submerged three-dimensional foil, and the full ship-strut-foil configuration. For the three-dimensional cases, the results shown correspond to a location near midspan of the foil, where results most closely approximate two-dimensional cases. For clarity, some of the pressure coefficients on the lower surface which are similar for cases already illustrated have been omitted.

The lift coefficients for the two-dimensional cases for all angles considered are shown in Table 1, along with the results for the three-dimensional cases.

DISCUSSION OF RESULTS

PRESSURE AND LIFT COEFFICIENTS

For the cases illustrating a change in pitch angle of the rigid ship-strut-foil configuration (Figures 2 through 6), the onset flow is essentially forced to be parallel to the foil and ship bottom when passing over the foil, regardless of the initial angle of approach. Thus, the force increase due to the change in pitch is relatively small (see Table 1), increasing about 60 percent over the range of 0° to $+5^{\circ}$ of pitch. In contrast, when the foil changes its position relative to the hull bottom (approximated by a flat plate), the force increases over 100 percent for a pitch change of $+2^{\circ}$ as shown in Table 1. This drastic change is born out in both the two-dimensional and three-dimensional cases. For the two-dimensional case, the C_p value for the $+5^{\circ}$ pitch angle has a minimum value of -3.2 which is off the scale shown in Figure 11. These large C_p changes in the flat plate cases are due to the cumulative effects of positive angle of foil pitch and the presence of the flat plate, both of which tend to lower the values of C_p on the upper surface.

Figure 7 shows the location and magnitude of $-C_{p_{max}}$ for the upper and lower surfaces along the span of the foil for various angles of onset flow. The chordwise location of $-C_{p_{max}}$ changes from 17.5 percent of chord to 22.5 percent of chord on the upper surface of the foil as the angle of onset flow changes from 0° to $+5^{\circ}$. The largest value of $-C_{p_{max}}$ occurs immediately inboard of the intersection of the foil and strut. The presence of the strut significantly increases the magnitudes of $-C_{p_{max}}$ at the strut relative to the value of $-C_{p_{max}}$ near midspan. This can be seen readily by comparison of C_p values in Figures 2 through 6 at foil span locations of y/s = 0.33 and 0.047.

Figure 10 illustrates further the magnitude of the effect of the strut on the foil. The local lift coefficient, C_{LL} , is markedly lower on the outboard side of the strut, and the smooth decrease in lift, seen in the cases without a strut, is gone. In addition, as the angle of pitch becomes increasingly negative, C_{LL} can be driven negative near the tip.

Figures 12 and 13 show very clearly the differences in pressure coefficient $C_{\rm p}$ between the various two-dimensional and three-dimensional cases. Both figures show that on the lower side, the values of $C_{\rm p}$ for the various

cases are in relatively good agreement. Also, in both figures, the two deeply submerged cases yield the lowest magnitudes of C_p on the upper surface. At the pitch angle, α , of 0° , Figure 12 shows that the values of $-C_{p_{max}}$ on the upper surface are within 10 percent of each other for the ship-strut-foil case and the two flat plate cases. For a pitch of $+2^{\circ}$, Figure 13 shows that the flat plate cases have significantly larger magnitudes of $-C_{p_{max}}$ than the ship-strut-foil case.

In conclusion, the strut intersection exerts a significant effect on the flow over the upper surface of the foil, while not significantly affecting flow on the lower surface. In addition, change in angle of pitch of the entire ship-strut-foil system does not affect the C_p on the foil nearly as much as rotation of the foil relative to the hull bottom. Rotation of the foil with respect to the ship's hull has a significant effect on the flow over the foil's surfaces as well as on the total applied forces on the foil.

PREDICTION OF CAVITATION SPEED

It is of interest to compute the maximum theoretical speed at which a smooth foil may operate free of cavitation, i.e., the cavitation inception speed. With the assumption that cavitation begins when the local static pressure is equal to the vapor pressure p_V , the cavitation inception speed $V_{\rm T}$ is given by

$$V_{I} = \sqrt{\frac{p_{V} - (p_{O} + \rho gh)}{\frac{1}{2} \rho C_{pm}}}$$
 [2]

where

 ${\rm p}_{\rm o}$ is the static pressure at the ocean surface, g is the acceleration of gravity h is the depth below the free surface, and ${\rm C}_{\rm pm}$ is the value of ${\rm C}_{\rm p}$.

For the ship-strut-foil case, the minimum value of C_p on the foil always occurs on its upper side, which is submerged approximately 17 feet (5.18 m) below the free surface. Also, the minimum value of C_p is approximately equal to -1.4 in the range $-5^\circ < \beta \le +5^\circ$ (See Figures 2-6). Assuming $\rho = 1.99$ 1b-sec 2 /ft 4 (1025.6 kg/m 3) for the density of seawater,

 $p_0 = 2160 \text{ lb/ft}^2 \text{ (5224 N/m}^2)$, and $p_V = 30 \text{ lb/ft}^2 \text{ (72.59 N/m}^2)$, Equation [2] yields the following value of V_T for the ship-strut-foil case

$$V_{T} = 48.1 \text{ ft/sec} = 28.5 \text{ knots}$$
 [3]

For cases where the foil is rotated relative to the ship bottom, considerably lower values of $V_{\rm I}$ may be expected for large positive values of α due to lower values of $C_{\rm pm}$. Figure 9 shows that the value of $C_{\rm pm}$ for the three-dimensional foil deeply submerged and pitched at $2^{\rm O}$ is -0.87 and the $C_{\rm pm}$ for the three-dimensional foil pitched at $2^{\rm O}$ beneath a flat plate is -1.49. The use of Equation [2] yields the following values for $V_{\rm I}$ for these cases:

 V_{I} for the deeply submerged foil = 60.98 ft/sec = 36.13 knots V_{I} for the foil beneath a flat plate = 46.60 ft/sec = 27.61 knots

The V_I for the deeply submerged foil is obviously a most optimistic estimate since the effect of the ship and strut is neglected entirely. A more conservative estimate can be obtained by using the two-dimensional results. Figure 11 shows that the values of C_{pm} for the two-dimensional foil pitched $+2^{\circ}$ and $+5^{\circ}$ are respectively -1.82 and -3.19. The use of Equation [2] yields the following values for V_{2D} , the cavitation inception speed for the two-dimensional foil

$$V_{2D} = 42.2 \text{ ft/sec} = 25.0 \text{ knots at } \alpha = +2^{\circ} \text{ (4a)}$$

 $V_{2D} = 31.8 \text{ ft/sec} = 18.9 \text{ knots at } \alpha = +5^{\circ} \text{ (4b)}$

It should be noted, however, that all of the above computed values of cavitation inception speed include no provision for manufacturing inaccuracies or roughness could degrade overall cavitation performance.

ACKNOWLE DGEMENT

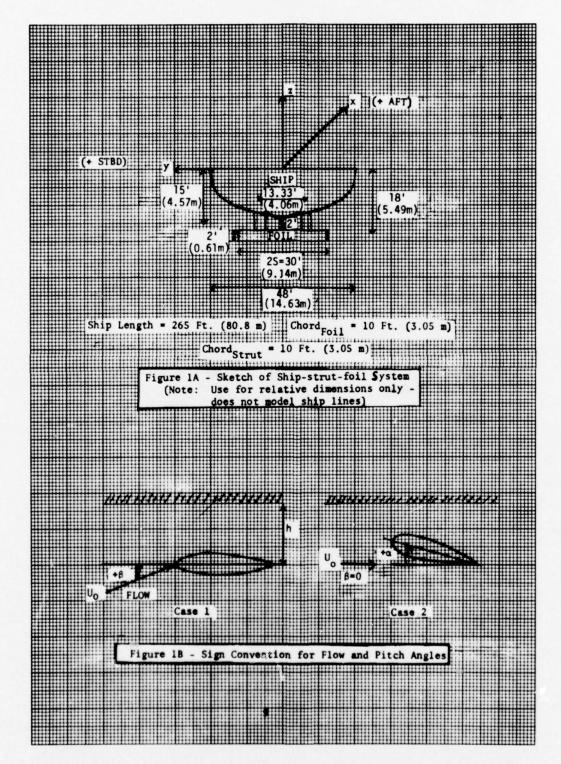
Acknowledgement is gratefully made to Mr. Charles W. Dawson for his assistance.

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- Smith, A.M.O., and J. Pierce, "Exact Solution of the Neumann Problem. Calculation of Non-Circulatory Plane and Axially Symmetric Flows About or Within Arbitrary Boundaries," Douglas Aircraft Company Report #ES-26988, April 1958.

TABLE 1 Summary of Lift Coefficients, $\mathbf{C}_{\mathbf{L}}$, For All Cases

3D Cases	Lift Coefficient C L Angle					
Case	Angle Varied	-5	-2	0	+2	+5
Ship-Foil-Strut Combination	β	0.055	0.10	0.14	0.17	0.22
Foil 2 Feet Below Flat Plate	α			0.19	0.44	
Foil Deeply Submerged	В			0	0.13	
2D Cases						
Foil 2 Feet Below Flat Plate	α	-0.65	-0.09	0.34	0.81	1.58
Foil Deeply Submerged	β	-0.681	-0.273	0	0.273	0.681



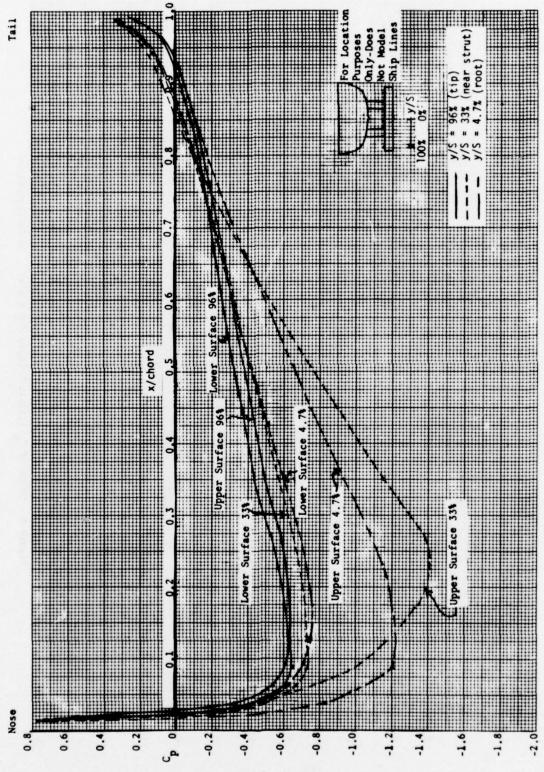
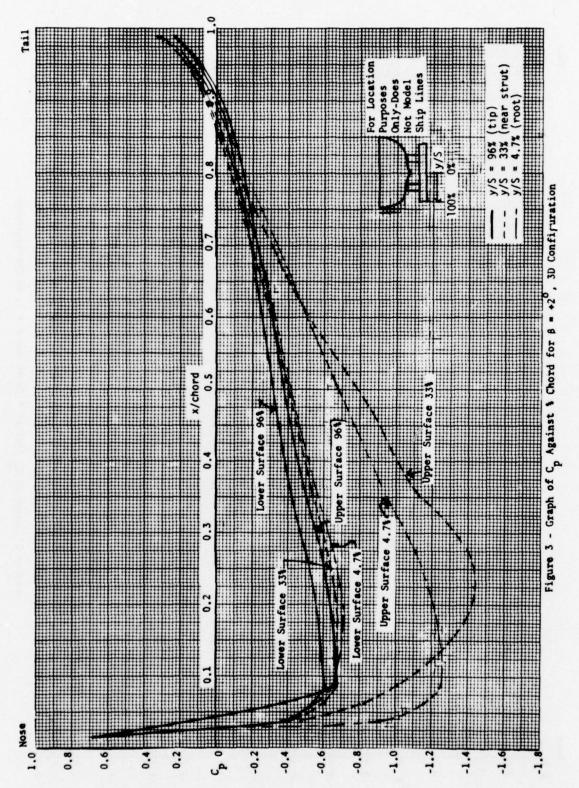
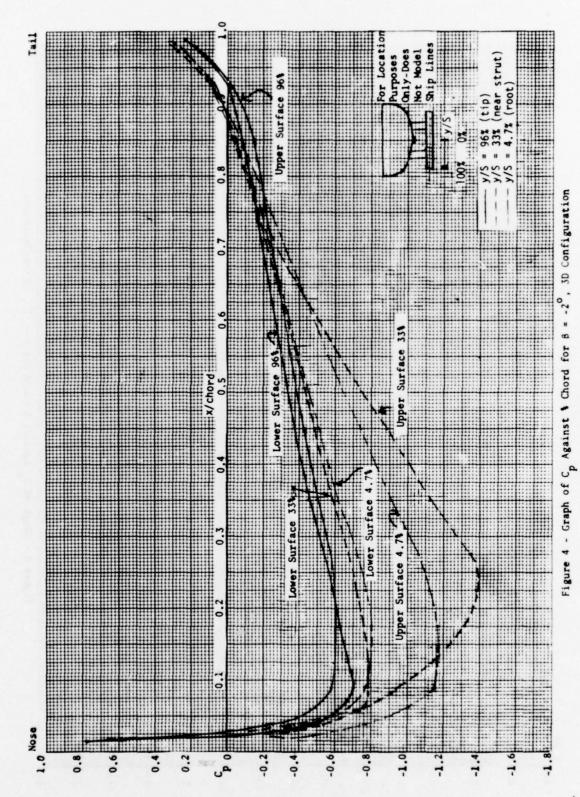
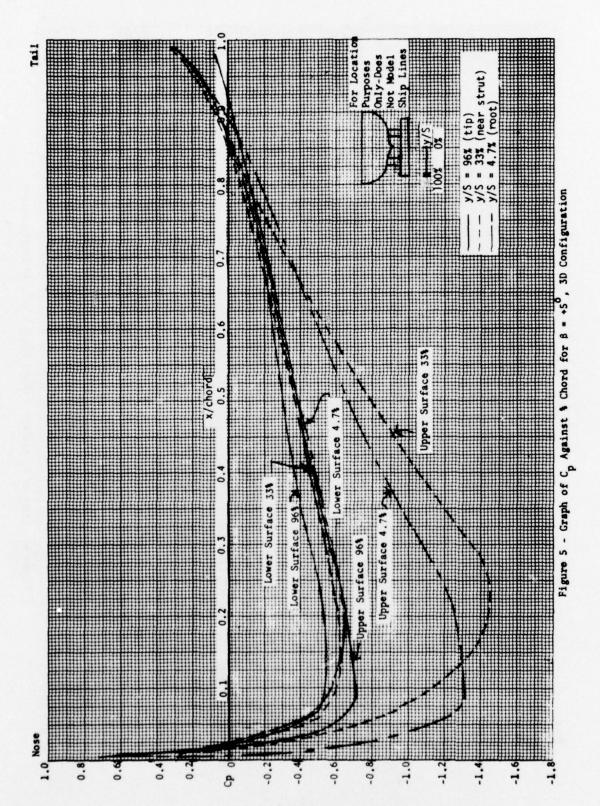
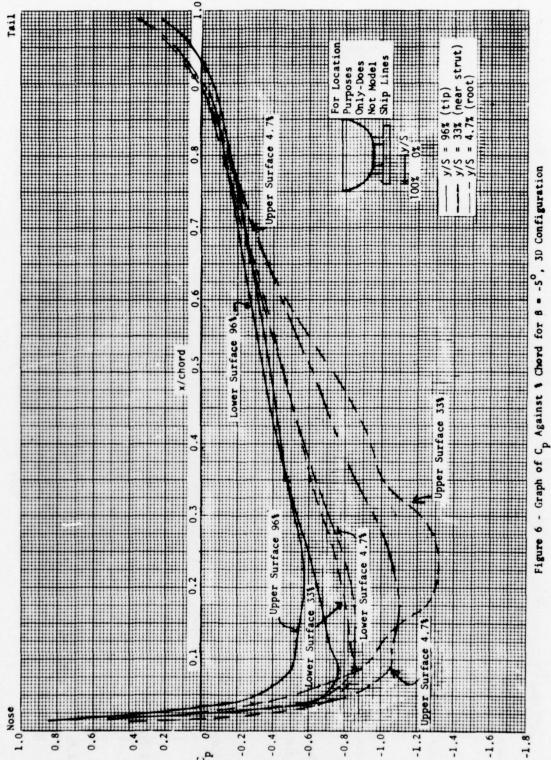


Figure 2 - Graph of C_p Against \$ Chord for β = 0, 3D Configuration









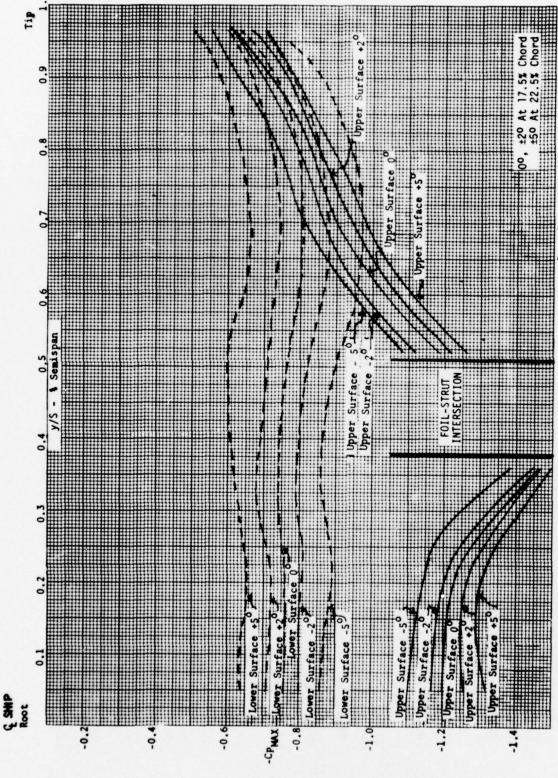
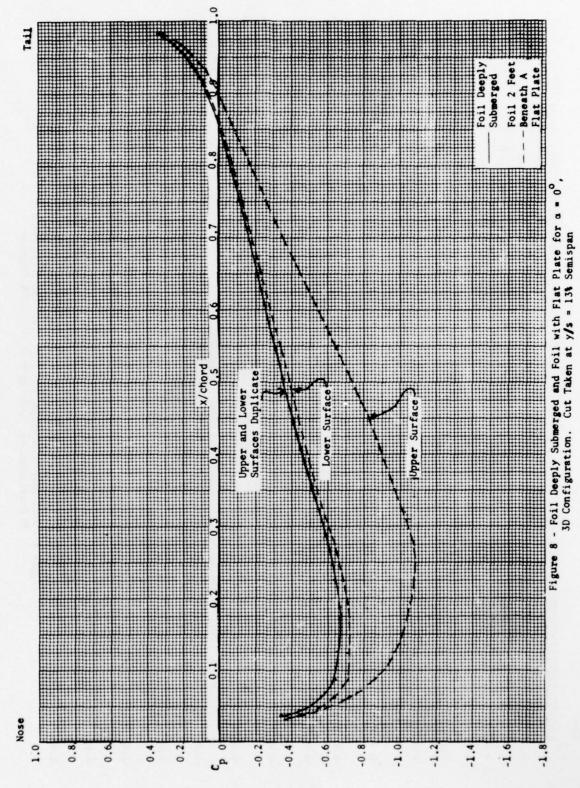


Figure 7 - Graph of -Cpmax Against & Semispan for β = 0, $\pm 2^{\circ}$, $\pm 5^{\circ}$, 3D Configuration



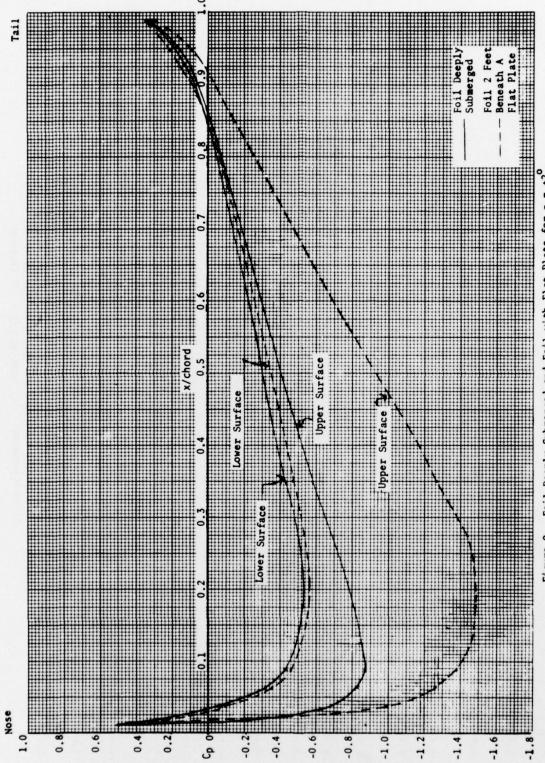
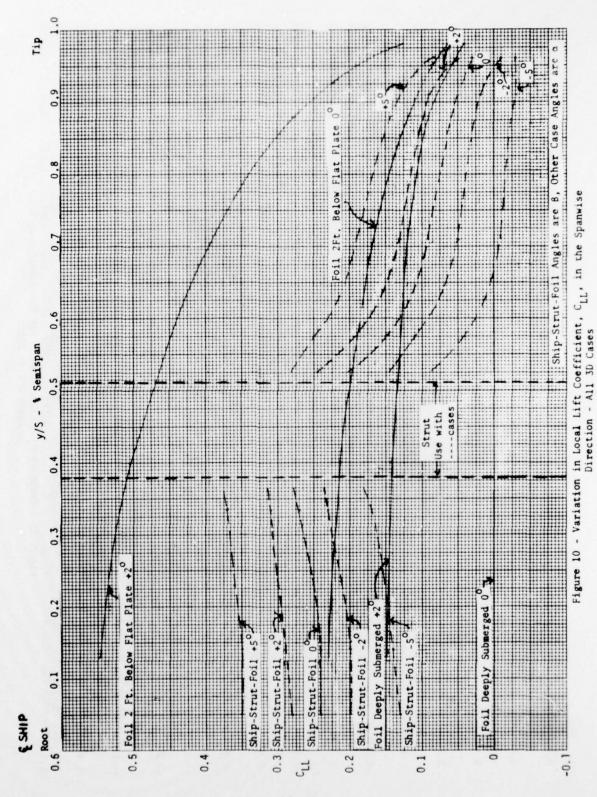
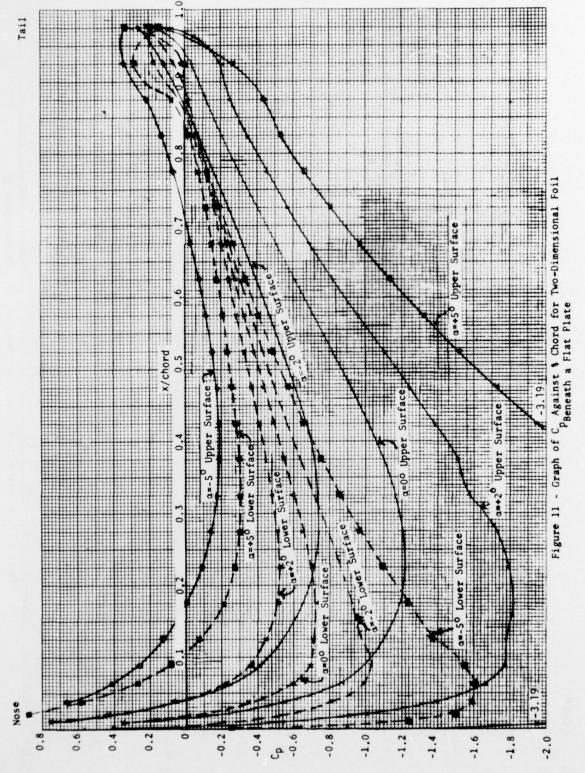
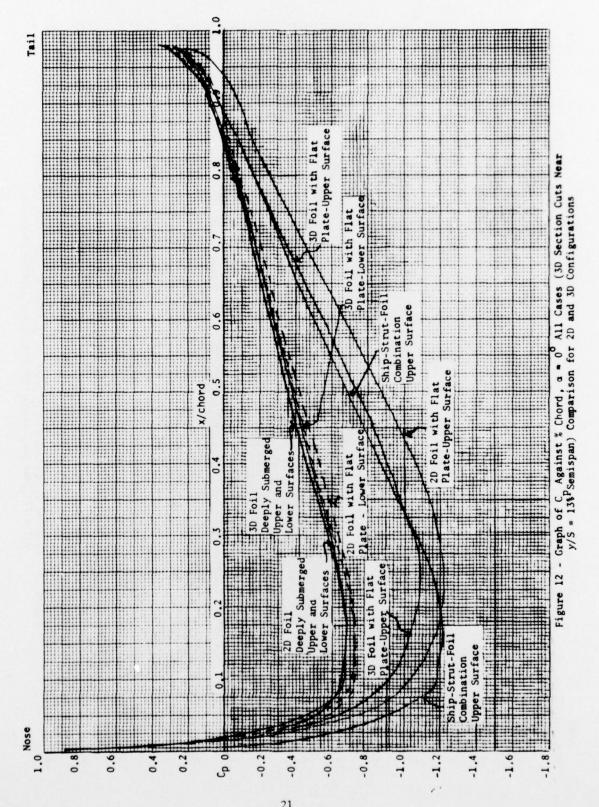
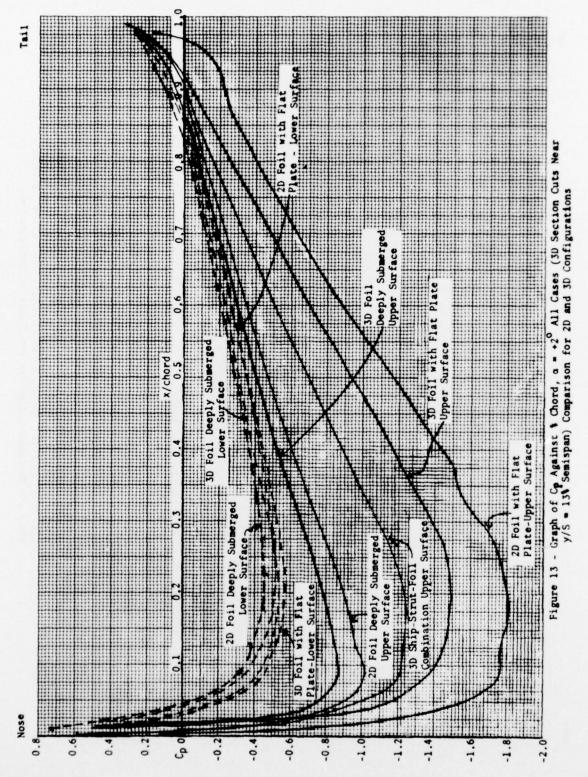


Figure 9 - Foil Deeply Submerged and Foil with Flat Plate for α = +2°, 3D Configuration. Cut Taken at y/s = 13% Semispan

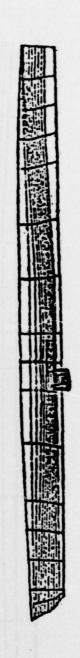




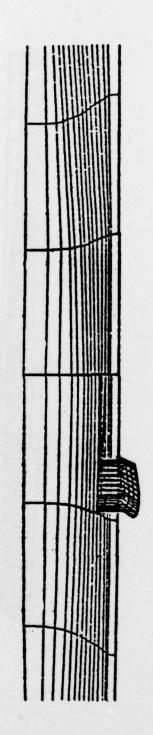




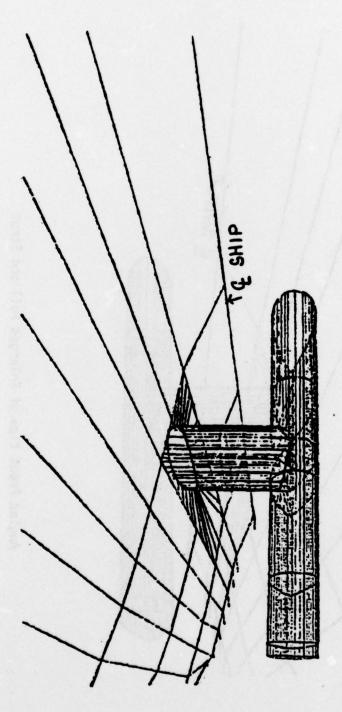
APPENDIX



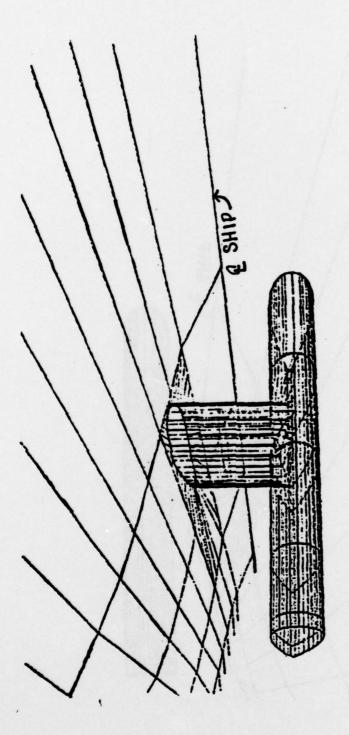
Graphics Representation of Entire Three-Dimensional Ship-Strut-Foil Configuration



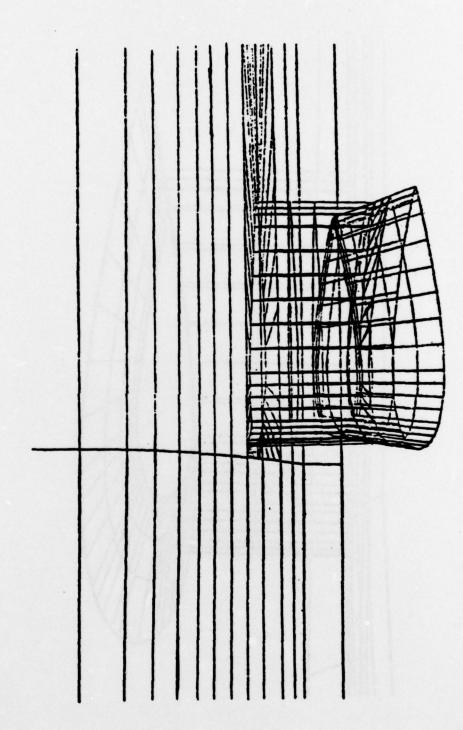
Graphics Representation of Foil and Strut Section - Enlarged



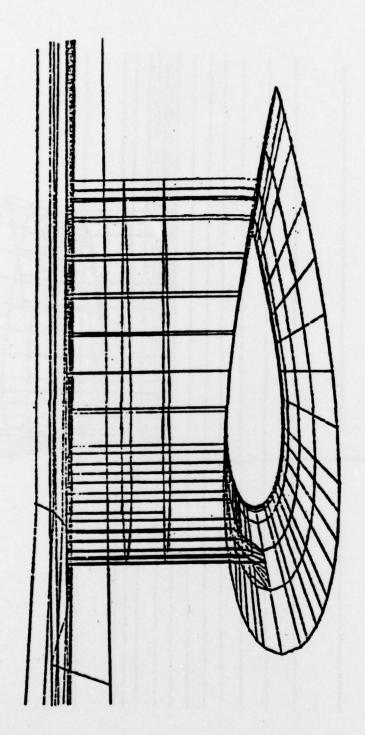
Front View of Enlarged Foil and Strut



Angled Front View of Enlarged Foil and Strut



Upper Surface of Foll - Detail



Side View of Strut-Foil Configuration

Enlarged View of Foil - Detail

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